

# Cinématique

① uniforme  $\Rightarrow v = \text{cte} \Rightarrow \tau = \frac{AB+BC+CD+DE}{v} = \frac{425\text{cm}}{25\text{cm/s}} = 17\text{sec} \Rightarrow \text{C}$

② circulaire uniforme  $\Rightarrow v = R\dot{\theta}$  et  $a_n = R\dot{\theta}^2 = \frac{R^2\dot{\theta}^2}{R} = \frac{v^2}{R} = \boxed{\approx 0,2\text{m}\cdot\text{s}^{-2}}$  (B)

$BC = \frac{2\pi R}{4} \Rightarrow R = \frac{2BC}{\pi} \Rightarrow a_n = \frac{(0,25)^2 \text{m}\cdot\text{s}^{-2}}{\frac{2 \times 0,50}{\pi} \text{m}} = \frac{\pi}{16} \text{m}\cdot\text{s}^{-2} \approx \frac{2 \times 157}{16}$

③  $a_2 = \frac{v^2 \times R}{R' \times R} = a_n \times \frac{R}{R'} = a_n \times \frac{\frac{2BC}{\pi}}{\frac{2CD}{\pi}} = a_n \times \frac{50}{75} = a_n \times \frac{2 \times 25}{3 \times 25} = \frac{2}{3} a_n = a_2$  (C)

normal vinge - brusque  
 $\Rightarrow a_1 > a_2$

④  $\begin{cases} x_a = \text{abs. de } M_a & x = 0 \text{ en D} \\ x_b = \text{abs de } M_b & x = 2\text{m en E} \end{cases}$

$x_a = \underset{25\text{cm}\cdot\text{s}^{-1}}{vt} \quad x_b = DE - \underset{50\text{cm}\cdot\text{s}^{-1}}{2vt}$

rencontre en  $x_a = x_b$   
i.e.  $vt_n = DE - 2vt_n \Rightarrow t_n = \frac{DE}{3v} = \frac{200\text{cm}}{75\text{cm}\cdot\text{s}^{-1}}$   
 $t_n = \frac{8 \times 25}{3 \times 25} \text{s} = 2,7 \text{sec.}$  (B)

⑤  $d_a = x_a(t_n) = vt_n = 25\text{cm/s} \times \frac{8}{3} \text{sec} = \frac{200}{3} = 67\text{cm}$  (B)

⑥ Soit  $t_0 = t_n - \frac{1}{5} \text{sec} = \frac{8}{3} - \frac{1}{5} = \frac{40-3}{15} = \frac{37}{15} \text{sec.}$

~~$x_a(t_0)$~~   $d = x_b(t_0) - x_a(t_0) = DE - 2vt_0 - vt_0 = DE - 3vt_0$

$d = 200\text{cm} - 3 \times 25 \times \frac{37}{15} \text{cm} = 200 - 5 \times 37 = 200 - 185 = 15\text{cm}$  (C)

# Dynamique

⑦  $\phi \text{ const} \Rightarrow E_m = \text{cte} \Rightarrow \frac{1}{2} m v_B^2 - \frac{1}{2} m v_A^2 = mg(z_A - z_B) = mg AB \sin \alpha$   
 $\Rightarrow v_B = \sqrt{2g AB \sin \alpha} = \sqrt{20 \times 40 \times \sin 30} = \sqrt{\frac{20 \times 40}{2}} = \sqrt{400} = 20 \text{ m} \cdot \text{s}^{-1}$   
 $\hookrightarrow$  (B) et (C)

⑧ Mvt circulaire mais pas uniforme  $\Rightarrow \begin{cases} \dot{\theta} \neq \text{cte} \\ \dot{r} = 0 \end{cases} \Rightarrow \begin{cases} \vec{v} = r \dot{\theta} \vec{e}_\theta \\ \vec{a} = r \ddot{\theta} \vec{e}_\theta - r \dot{\theta}^2 \vec{e}_r \end{cases}$

~~passer~~ À cet endroit  $\vec{R} \propto \vec{P}$  donc  $m \vec{a} = \vec{R} + \vec{P}$  suivant  $\vec{e}_\theta$  donc  $\vec{a}_c \cdot \vec{e}_\theta = 0$  (B)

en C l'intensité passe par un max donc  $\dot{\theta}$  est max donc  $\ddot{\theta} = 0$  (D) juste

⑨  $z_B = z_0 - b \cos \alpha = b \cos \beta - b \cos \alpha \Rightarrow E_p(B) = mgb(\cos \beta - \cos \alpha)$  (D)

⑩  $E_m = \text{cte} \Rightarrow \frac{1}{2} m v_0^2 - \frac{1}{2} m v_B^2 = mg z_B \Rightarrow v_0 = \sqrt{v_B^2 + \frac{2E_p(B)}{m}}$  (C)

$v_0 < v_B$  car + haut

⑪  $v_0 = v_0 \leftarrow$  j'imagine faute de type.  
 $\begin{matrix} \uparrow & \uparrow \\ z_{\text{no}} & \text{"eau"} \end{matrix}$

$a_z = -g \Rightarrow v_z = -gt + v_0 \sin \beta \rightarrow t_{\text{max}} = \frac{v_0 \sin \beta}{g}$  (C)  
 $\hookrightarrow z(t) = -\frac{1}{2} g t^2 + v_0 \sin \beta t \Rightarrow z_F = -\frac{v_0^2 \sin^2 \beta}{2g} + \frac{v_0^2 \sin^2 \beta}{g}$   
 $z_F = +\frac{v_0^2 \sin^2 \beta}{2g} = \frac{v_0^2 \times \left(\frac{\sqrt{2}}{2}\right)^2}{2g} = \frac{v_0^2}{4g}$

⑫  $\theta$  angle p.n à la verticale O'C  
 $\vec{OM} = b \vec{e}_r; \vec{v} = b \dot{\theta} \vec{e}_\theta; \vec{a} = b \ddot{\theta} \vec{e}_\theta - b \dot{\theta}^2 \vec{e}_r = \frac{\vec{P} + \vec{R}}{m} = +g \cos \theta \vec{e}_r - g \sin \theta \vec{e}_\theta - R \vec{e}_r$   
 $\hookrightarrow$  suivant  $\vec{e}_\theta$  :  $b \ddot{\theta} = -g \sin \theta \approx -g \theta$  car  $\theta \ll 1$



d'où  $\ddot{\theta} + \frac{g}{b} \theta = 0 \Rightarrow T = \frac{2\pi}{\omega} = 2\pi \sqrt{\frac{b}{g}} = 2 \times 3,14 \times \sqrt{2}$   
 $= 6,3 \times 1,4 = 8,82 \approx 9 \text{ sec}$  (C)

## Electrocinétique - transitoires

(13)  $i_{\text{Bob}}$  continue  $\Rightarrow i_L(0^+) = i_R(0^+) = 0 \Rightarrow u_R(0^+) = R i_R(0^+) = 0$  (A)  
 $\Rightarrow u_L(0^+) = E$  (D)

(14)  $E = L \frac{di_L}{dt} + R i_L \Rightarrow \frac{di_L}{dt} + \frac{R}{L} i_L = \frac{E}{L} \Rightarrow i_L = \frac{E}{R} [1 - \exp(-\frac{R}{L} t)]$  (A)

$u_L = L \frac{di_L}{dt} = \frac{EL}{R} \left( 0 + \frac{R}{L} \exp(-\frac{R}{L} t) \right) = E \exp(-\frac{R}{L} t)$  (D)

(15)  $i_L(0^+) = i_L(0^-) = \lim_{t \rightarrow +\infty} \left( \frac{E}{R} (1 - e^{-\frac{R}{L} t}) \right) = \frac{E}{R}$  (B)

$u_R(0^+) = u_C(0) = \frac{E}{3} \Rightarrow i_R(0^+) = \frac{E}{3R}$  (D)

(16)  $i_C(0^+) = i_L(0^+) - i_R(0^+) = \frac{3E}{3R} - \frac{E}{3R} = \frac{2E}{3R}$  (B)

$u_L(0^+) = E - \frac{E}{3} = \frac{2E}{3}$  (C)

(17)  $E = L \frac{di_L}{dt} + R i_R = L \frac{di_L}{dt} + R (i_L - i_C) = L \frac{di_L}{dt} + R i_L - RC \frac{du_C}{dt}$

$E = L \frac{di_L}{dt} + R i_L - RC \left( \frac{dE}{dt} - \frac{du_L}{dt} \right) = L \frac{di_L}{dt} + R i_L - RC \left( 0 - L \frac{d^2 i_L}{dt^2} \right)$

$E = +RLC \frac{d^2 i_L}{dt^2} + L \frac{di_L}{dt} + R i_L \Leftrightarrow \frac{E}{RLC} = \frac{d^2 i_L}{dt^2} + \frac{1}{RC} \frac{di_L}{dt} + \frac{1}{LC} i_L$

$\omega_0 = \frac{1}{\sqrt{LC}}$  et  $\tau_c = RC$

$\frac{E}{RLC} = \frac{d^2 i_L}{dt^2} + \frac{1}{RC} \frac{di_L}{dt} + \frac{1}{LC} i_L$

(C) (A)

(18)  $\frac{E}{RLC} = \omega_0^2 i_\infty = \frac{i_\infty}{LC} \Rightarrow i_\infty = \frac{E}{R}$  (B)

$t \rightarrow \infty \Rightarrow \text{Bob} = \text{Fil} \Rightarrow u_C(\infty) = E \Rightarrow \mathcal{E}_C = \frac{1}{2} C E^2$  (D)

# Régime sinusoïdal forcé

19

$$\frac{U_{C,m}}{U_{e,m}} = \frac{1/j\omega C}{1/j\omega C + j\omega L + R + r_L + r_i} = \frac{1}{1 + j(R_{TOT} \omega - L\omega^2)}$$

$$H_0 = 1 \quad (A)$$

$$\omega_0 = \frac{1}{\sqrt{LC}} \quad (D)$$

$$\frac{U_{C,m}}{U_{e,m}} = \frac{1}{1 + j \frac{R_{TOT}}{L} \times \frac{\omega}{\omega_0} - \left(\frac{\omega}{\omega_0}\right)^2} = \frac{1}{1 + j \frac{R_{TOT}}{L\omega_0} \omega - \omega^2}$$

$$= \frac{1}{Q} \rightarrow Q = \frac{L\omega_0}{R_{TOT}} = \frac{1}{R + r_L + r_i} \sqrt{\frac{L}{C}}$$

20 → (B)

21

pdt

$$\begin{cases} U_{C,m} = U_{e,m} \times \frac{1/j\omega C}{1/j\omega C + j\omega L + R + r_L} \\ U_{L,m} = e_m \times \frac{1/j\omega C}{1/j\omega C + j\omega L + R + r_L + r_i} \end{cases}$$

$$\left(\frac{U_{C,m}}{U_{e,m}}\right)^{-1} = \frac{1/j\omega C + j\omega L + R + r_L + r_i}{1/j\omega C + j\omega L + R + r_L}$$

$$\text{or } \omega = \omega_0 = \frac{1}{\sqrt{LC}} \text{ donc}$$

$$\frac{1}{j\omega_0 C} = -j\sqrt{\frac{L}{C}} \text{ et } jL\omega_0 = +j\sqrt{\frac{L}{C}}$$

$$\Rightarrow \frac{1}{j\omega_0 C} = -jL\omega_0 \text{ et } \dots$$

(C)

22

$$i_m = \frac{e_m}{jL\omega_0 + \frac{1}{j\omega_0 C} + R + r_L + r_i} \Rightarrow (B)$$

$$U_{C,m} = \frac{1}{1 + j/Q} \times U_{e,m} = -jQ \times U_{e,m} = -jQH_0 U_{e,m}$$
 ↑  $1 + j/Q = 1 - j/Q$   
 question 19

23

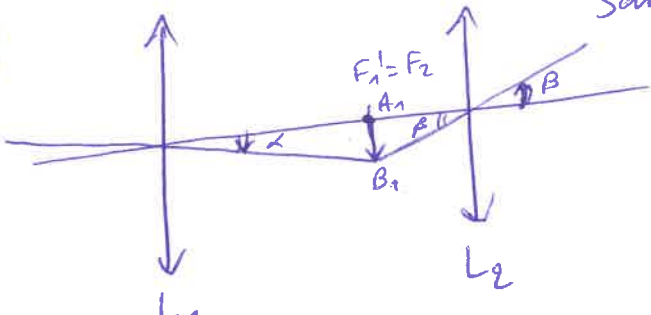
$$U_{L,m} = (jL\omega_0 + r_L) i_m = jL\omega_0 \frac{e_m}{R + r_L + r_i} + \frac{r_L e_m}{R + r_L + r_i}$$

$$U_{L,m} = \left( \frac{r_L}{R + r_L + r_i} + j \dots \right) e_m \rightarrow b = \frac{1}{R + r_L + r_i} \times \sqrt{\frac{L}{C}}$$

$$\hookrightarrow a \Rightarrow (A) \quad (24) \rightarrow (C)$$

# Optique

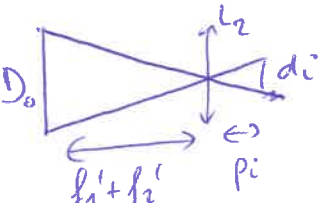
(25) (D) en accommodant ~~non~~

(26)  sans accommoder  $\Rightarrow F_1' = F_2'$

$$G_a = \frac{\beta}{\alpha} \approx \frac{\tan \beta}{\tan \alpha} \approx -\frac{\frac{A_1 B_1}{f_2'}}{\frac{A_1 B_1}{f_1'}} = -\frac{f_1'}{f_2'} \quad (B)$$

(27)  $\frac{1}{p_i} + \frac{1}{f_1' + f_2'} = \frac{1}{f_2'} \Rightarrow \frac{1}{p_i} = \frac{-1}{f_1' + f_2'} + \frac{1}{f_2'} = \frac{f_1'}{f_2'(f_1' + f_2')} \Rightarrow p_i = \frac{f_2'(f_1' + f_2')}{f_1'}$

$$p_i = f_2' \left( 1 + \frac{f_2'}{f_1'} \right) \quad (D)$$

  $d_i = \frac{D_0 \times p_i}{f_1' + f_2'} = D_0 \times \frac{f_2'}{f_1'} \quad (A)$

(28) On veut  $\overline{O_2 A'} = -dm$  avec  $\overline{O_2 A_1} = \overline{O_2 F_1'} = -(f_2' - x)$  donc :

↑  
image finale

$$-\frac{1}{dm} + \frac{1}{f_2' - x} = \frac{1}{f_2'} \Rightarrow \frac{1}{f_2' - x} = \frac{dm + f_2'}{f_2' dm} \Rightarrow x = f_2' - \frac{f_2' dm}{dm + f_2'}$$

$$x = \frac{f_2' dm + f_2'^2 - f_2' dm}{dm + f_2'} \Rightarrow (A)$$

(29) Newton :  $\overline{F_1' A} \times \overline{F_1' B} = -f_1'^2$     afocal  $\Rightarrow F_1' = F_2'$

(A)  $\left\{ \begin{array}{l} \overline{F_2' B} \times \overline{F_2' C} = -f_2'^2 \\ \overline{F_1' B} \times \overline{F_2' C} = -f_2'^2 \end{array} \right. \quad (D)$

(30) 
$$\begin{cases} -D_{\text{min}} \times \overline{F_1' B} = -f_1'^2 \\ -\overline{F_1' B} \times (f_2' + dm) = -f_2'^2 \end{cases} \quad (C)$$

$$D_{\text{min}} = \left( \frac{f_1'}{f_2'} \right)^2 \times (f_2' + dm)$$

# Thermodynamique

(31)  $\frac{V_1}{T_1} = \frac{nR}{P_1} = \frac{nR}{P_2} = \frac{V_2}{T_2} \Rightarrow V_2 = \frac{T_2}{T_1} V_1$  ; de +  $V_1 + V_2 = 2V_0 = V_t$   
 can  $P_1 = P_2$  eq. méca      donc  $V_1 \left(1 + \frac{T_2}{T_1}\right) = V_t \Rightarrow V_1 = \frac{T_1}{T_1 + T_2} V_t$

eq. thermiq gaz 2  $\Rightarrow T_2 = T_0 \Rightarrow V_1 = \frac{T_1}{T_1 + T_0} V_t$  (B)

$V_2 = V_t - V_1 = V_t \left(1 - \frac{T_1}{T_1 + T_0}\right) = V_t \left(\frac{T_1 + T_0 - T_1}{T_1 + T_0}\right) = \frac{T_0}{T_1 + T_0} V_t$  (D)

(32)  $\rightarrow$  (A) et  $P_1 = \frac{nRT_1}{V_1} = \frac{nRT_1(T_1 + T_0)}{T_1 V_t} = \frac{nR(T_0 + T_1)}{V_t}$  (B)

(33)  $\Delta U = \Delta U_1 + \Delta U_2 = nC_{vm}(T_1 - T_0 + T_0 - T_0) = \frac{nR}{\gamma - 1}(T_1 - T_0)$  (B)  
 $\Delta U_2 = 0$  can GP et  $T_i = T_f$

MAVER:  $\begin{cases} C_{pm} - C_{vm} = R \\ \frac{C_{pm}}{C_{vm}} = \gamma \end{cases} \Rightarrow C_{vm}(\gamma - 1) = R$

(34)  $\Delta U_2 = 0 = W_2 + Q_2 \Rightarrow W_2 = -Q_2$  (D)  
 rév  $\Rightarrow$  eq. thermique constant

$W_2 = \int_{V_0}^{V_2} -pdV = - \int_{V_0}^{V_2} \frac{nRT_0}{V} dV = -nRT_0 \ln \frac{V_2}{V_0} = nRT_0 \ln \frac{V_0}{V_2} = nRT_0 \ln \frac{V_t}{2V_2}$   
 $V_0$   $C_p = p_{ext}$  quasistatique

$W_2 = nRT_0 \ln \frac{V_t(T_0 + T_1)}{2T_0 V_t}$  (A)

(35)  $W_1$  reçu par gaz 1 =  $-W_2 \Rightarrow \Delta U_1 = Q_1 - W_2 \Rightarrow Q_1 = \Delta U_1 + W_2$  or  
 on a montré que  $\Delta U_2 = 0$  donc  $Q_1 = \Delta U + W_2$  (B).

Pélec =  $Ri^2 \Rightarrow Q_1 = Ri^2 \tau$  (D)

(36)  $S_2^{(n)} = \frac{Q_2}{T_0} = \frac{-W_2}{T_0} = -nR \ln \left(\frac{T_0 + T_1}{2T_0}\right) = nR \ln \left(\frac{2T_0}{T_0 + T_1}\right)$  (A)

$\Delta S_2 = S_2^{(n)} + S_2^{(c)} \Rightarrow S_2^{(c)} = \Delta S_2 - S_2^{(n)} = nR \ln \left(\frac{T_0}{T_0 + T_1} \times \frac{V_t}{V_0}\right) - nR \ln \left(\frac{2T_0}{T_0 + T_1}\right)$

$S_2^{(c)} = 0$  (D)  $\frac{V_t}{V_0} = \frac{2T_0}{T_0 + T_1}$   
 = 0 normal réversible.